

The chaotic heat equation

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Abstract. The one dimensional heat equation

$$\frac{\partial u}{\partial t} = a \frac{\partial^2 u}{\partial x^2} + b \frac{\partial u}{\partial x} + cu$$

for $x \in \mathbb{R}$, $t \geq 0$ is governed by a chaotic semigroup for certain values of the coefficients $(a, b, c) \in \mathbb{R}^3$ on certain weighted supremum norm spaces.

1 Introduction

In 1997, G. Herzog [1] proved the universality of the heat equation in an exponentially weighted space. Concerning the universal families K. Grosse-Erdmann published an expository paper [2] in which we can find the following definition.

Definition 1. Let X and Y be topological spaces and $T_t : X \rightarrow Y$ ($t \in I$) continuous mappings. Then an element $x \in X$ is called *universal* (for the family $\{T_t\}_{t \in I}$) if the set $\{T_t x : t \in I\}$ is dense in Y . The set of universal elements is denoted by $\mathcal{U} = \mathcal{U}(T_t)$. The family $\{T_t\}_{t \in I}$ is called *universal* if it has a universal element.

The set \mathcal{U} of universal elements may be expected to be huge, in fact residual in the sense of Baire categories. In [1] G. Herzog defined for $\alpha > 0$, the separable Banach space

$$E_\alpha := \{\varphi \in C(\mathbb{R}, \mathbb{R}) : \lim_{x \rightarrow \pm\infty} e^{-\alpha|x|} \varphi(x) = 0\}$$

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equipped with the norm

$$\|\varphi\|_{E_\alpha} = \sup_{x \in \mathbb{R}} e^{-\alpha|x|} |\varphi(x)|$$

and by taking $T_t : E_\alpha \rightarrow C(\mathbb{R}, \mathbb{R})$,

$$[T_t \varphi](x) = \frac{1}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} \exp\left(-\frac{(x-s)^2}{4t}\right) \varphi(s) ds, \quad (1.1)$$

he showed that the set $\mathcal{U} := \{\varphi \in E_\alpha : \overline{\{T_n \varphi : n \in \mathbb{N}\}} = C(\mathbb{R}, \mathbb{R})\}$ is a residual subset of E_α for some $\alpha > 0$.

In operator theory, when one studies the iterates $\{T^n\}_{n \in \mathbb{N}_0}$ of a (continuous linear) operator T , universal elements are usually called *hypercyclic*. More generally a (C_0) semigroup $T = \{T(t)\}_{t \geq 0}$ on a separable Banach space X is called *hypercyclic*, if there exists an element $x \in X$ such that its orbit $\{T(t)x : t \geq 0\}$ is dense in X .

The idea of universality is well known in topological dynamics under the name of *topological transitivity*. This definition has been translated for a (C_0) semigroup in the following manner.

Definition 2. Let $T = \{T(t)\}_{t \geq 0}$ be a (C_0) semigroup on a separable Banach space X . T is called *chaotic* if it is hypercyclic and there are vectors $y \in X$ such that y has a periodic orbit, i.e., there exists $\tau = \tau(y) > 0$ such that $T(\tau)y = y$; and the set of such y 's is dense in X .

Recently F. Astengo and B. Di Blasio, in a nice paper [3], proved that if $A(x)$ defined in [3, p. 50] is the density of an absolutely continuous measure μ on the half line $(0, \infty)$, if $L^p = L^p((0, \infty), \mu)$ with $d\mu(x) = A(x)dx$, is the corresponding Lebesgue space with $2 < p < \infty$, and if $\Delta_p = d^2/dx^2$ is the Laplacian on L^p with maximal domain, then there is a constant $B(p)$ such that if $\beta > B(p)$, then $\Delta_p + \beta I$ generates a (C_0) chaotic semigroup on L^p . In this paper we prove a similar result in spaces of continuous functions.

In [4, 5] we have proved that the Black-Scholes semigroup is chaotic in certain Banach spaces which will be define later. The Black-Scholes (or Black-Merton-Scholes) equation is

$$\frac{\partial u}{\partial t} = \frac{\sigma^2}{2} y^2 \frac{\partial^2 u}{\partial y^2} + ry \frac{\partial u}{\partial y} - ru, \quad (1.2)$$

where $u = u(y, t)$, $y > 0$, $t \geq 0$, and r, σ are positive constants representing the interest rate and the volatility respectively. In a series of papers in the period 1973-1976, F. Black, R. Merton and M. Scholes showed that the fair value of a stock option u satisfies this equation and developed a significant economic model around it. Here y is the current value of the asset. It was surprising to deduce a deterministic model for a study which started with an Ito stochastic differential equation. This equation had a profound effect on the development of financial mathematics and it led to Merton and Scholes receiving the 1998 Nobel prize in Economics; by that time Black had passed away.

The initial condition this model requires is

$$u(y, 0) = \max\{0, y - p\}$$

for a fixed given $p > 0$. Thus a rigorous study of the semigroup governing (1.2) must take place on a function space which includes functions which grows at least linearly at infinity. On this it is natural of use supremum norm spaces with weights, namely, $X = \{f \in C(0, \infty) : f\rho \in C_0(0, \infty)\}$ where ρ is a continuous weight function on $(0, \infty)$ which grows faster than linearly at infinity. Here $C_0(0, \infty)$ consists of those continuous functions h on $(0, \infty)$ such that $\lim_{x \rightarrow L} h(x) = 0$ for $L \in \{0, \infty\}$.

The idea of *chaos* is most easily understood using the simple example of J. von Neumann in which the iterative process $x_{n+1} = 4x_n(1 - x_n)$ generates a chaotic sequence. In his book [6], R. Devaney defines the notion of chaos for a dynamical system by basing itself on this example. Let

$$F : [0, 1] \mapsto [0, 1]$$

be defined by

$$F(x) = 4x(1 - x).$$

The map $E : x \rightarrow \theta$, where $x = \sin^2 \theta$, is a diffeomorphism from $[0, 1]$ to $[0, \pi/2]$. Since $F(x) = 4\sin^2 \theta(1 - \sin^2 \theta) = \sin^2(2\theta)$, we deduce that the n -fold composition F^n satisfies

$$F^n(x) = \sin^2(2^n \theta)$$

for all $n = 1, 2, \dots$. Then the orbits $\{F^n(x)\}_{n=0}^\infty$ are dense in $[0, 1]$ whenever $\theta/2\pi$ is irrational, and they are periodic whenever $\theta/2\pi$ is rational. Thus F is a chaotic (nonlinear) map.

2 The chaotic Black-Scholes semigroup

Let $s, \tau \geq 0$ and define

$$Y_{s,\tau} := \left\{ f \in C(0, \infty) : \frac{f(x)}{(1+x^s)(1+x^{-\tau})} \in C_0(0, \infty) \right\}$$

with the norm

$$\|u\|_{s,\tau} = \sup_{x \geq 0} \left| \frac{u(x)}{(1+x^s)(1+x^{-\tau})} \right| < \infty.$$

The above space is a complex separable Banach space.

Theorem 1. ([4, 5]) *Let $s \geq 1, \tau \geq 0$. Let σ, r be given positive constants. Let $S = \{S(t) : t \geq 0\}$ be the semigroup governing the Black-Scholes equation (1.2) on $Y_{s,\tau}$. Then S is a (C_0) semigroup. Moreover, S is chaotic if $s > 1$. The same conclusions hold for the restriction of S to*

$$Y_{s,\tau}^{\mathbb{R}} = \{f \in Y_{s,\tau} : f \text{ is a real function}\}.$$

Note that (1.2) determines a chaotic semigroup on any of the spaces

$$\{Y_{s,\tau} : \tau \geq 0, s > 1\}.$$

3 The chaotic heat equation

The exponential function $y = e^x$ is a Lie group diffeomorphism from $(\mathbb{R}, +)$ to $((0, \infty), \cdot)$. This induces a map

$$E : f \mapsto g,$$

mapping functions on \mathbb{R} to function on $(0, \infty)$, defined by

$$(Ef)(y) = f(\log y) = g(y)$$

for $y \in (0, \infty)$ and $x = \log y \in \mathbb{R}$ so that $y = e^x$. Let w be a positive continuous function on \mathbb{R} . Let X_w be the weighted sup norm space

$$X_w = \{f \in C(\mathbb{R}) : \frac{f}{w} \in C_0(\mathbb{R})\}$$

with $\|f\|_{X_w} = \sup_{x \in \mathbb{R}} |f(x)|/w(x)$. Then E is an isometric isomorphism from X_w to $Y_{\tilde{w}}$, where

$$Y_{\tilde{w}} = \{g \in C(0, \infty) : \frac{g}{\tilde{w}} \in C_0(0, \infty)\},$$

equipped with $\|g\|_{Y_{\tilde{w}}} = \sup_{y > 0} |g(y)|/\tilde{w}(y)$, and $\tilde{w} = Ew$.

Let $Y_{s,\tau}$ be as in Section 2. Define

$$X_{s,\tau} = X_w,$$

where $w(x) = (1 + e^{-\tau x})(1 + e^{sx})$, for $x \in \mathbb{R}$. Then

$$\tilde{w}(y) = (Ew)(y) = (1 + y^{-\tau})(1 + y^s),$$

and $Y_{s,\tau} = Y_{\tilde{w}}$.

Now we work forward the goal of interpreting the result of [4] for functions on $(0, \infty)$ in the context of the functions on \mathbb{R} . The translation group T_a for functions on \mathbb{R} with parameter a [resp., S_a for functions on $(0, \infty)$] is given by

$$[T_a(t)f](x) = f(x + at) \quad \text{for } x \in \mathbb{R}, t \in \mathbb{R}$$

[resp.

$$[S_a(t)g](y) = g(e^{at}y) \quad \text{for } y > 0, t \in \mathbb{R}].$$

Here a is a fixed positive constant. The generator is, formally,

$$A_a = T'_a(0) = a \frac{d}{dx} \quad [\text{resp. } B_a = S'_a(0) = ay \frac{d}{dy}].$$

Then T_a [resp. S_a] is a (C_0) isometric group on $C_0(\mathbb{R})$ [resp. on $C_0(0, \infty)$]. In [4] it was shown that S_a is a (C_0) group on $Y_{s,\tau}$ for $s > 1, \tau \geq 0$. Since E is an isometric isomorphism from $X_{s,\tau}$ to $Y_{s,\tau}$, T_a is a (C_0) group on $X_{s,\tau}$ for the same values of s and τ .

Now suppose G generates a (C_0) group $\{e^{tG} : t \geq 0\}$ on a Banach space Z . Then G^2 generates a (C_0) semigroup on Z , analytic in the right half plane, given by

$$e^{tG^2} h = \frac{1}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} e^{-zG} h dz,$$

for $h \in Z$ and $\text{Re}(t) > 0$. This result is due to N. P. Romanoff (see [7] and [8, p. 120]). Consequently,

$$e^{tA_a^2} f(x) = \frac{1}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} f(x - az) dz, \quad (3.1)$$

$$e^{tB_a^2} g(y) = \frac{1}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} g(e^{-az}y) dz. \quad (3.2)$$

Since $A_a^2 = a^2 \frac{d^2}{dx^2}$ and $B_a^2 = a^2 y^2 \frac{d^2}{dy^2} + a^2 y \frac{d}{dy}$, the Black-Scholes equation (1.2) is

$$\begin{aligned}\frac{du}{dt} &= B_{\sigma/\sqrt{2}}^2 u + \left(r - \frac{\sigma^2}{2}\right) B_1 u - r I u, \\ &= \sum_{j=1}^3 G_j u.\end{aligned}\quad (3.3)$$

The three generators G_1, G_2, G_3 all commute, thus $e^{t\sum_{j=1}^3 G_j} = \prod_{j=1}^3 e^{tG_j}$, and the three factors in the product can be written in any order. Specifically, we have

$$G_1 = B_{\sigma/\sqrt{2}}^2, \quad G_2 = \left(r - \frac{\sigma^2}{2}\right) B_1, \quad G_3 = -rI.$$

Consequently the unique solution of (3.3) with initial condition $u(0) = g$ is given by

$$\begin{aligned}u(t, y) &= e^{tG_3} e^{tG_1} e^{tG_2} g \\ &= \frac{e^{-rt}}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} [e^{tG_2} g] \left(e^{-\frac{\sigma}{\sqrt{2}}z} y\right) dz \\ &= \frac{e^{-rt}}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} g \left(e^{-\frac{\sigma}{\sqrt{2}}z} e^{(r-\frac{\sigma^2}{2})t} y\right) dz.\end{aligned}$$

This formula was derived in a slightly more complicated but equivalent form in [9] and [10, p. 162].

The heat equation corresponding to (3.3) is

$$\begin{aligned}\frac{du}{dt} &= A_{\sigma/\sqrt{2}}^2 u + \left(r - \frac{\sigma^2}{2}\right) A_1 u - r I u, \\ &= \sum_{j=1}^3 H_j u.\end{aligned}\quad (3.4)$$

Using this and (3.1), we conclude that the unique solution (in $X_{s,\tau}$) of (3.4) is given by

$$\begin{aligned}u(t, y) &= \frac{e^{-rt}}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} \left[e^{t(r-\frac{\sigma^2}{2})} f \right] (x-z) dz \\ &= \frac{e^{-rt}}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-z^2/4t} f(x-z - (r-\frac{\sigma^2}{2})t) dz,\end{aligned}$$

provided $u(0, x) = f(x)$, $f \in X_{s,\tau}$. We rewrite the initial value problem as

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{\sigma^2}{2} \frac{\partial^2 u}{\partial x^2} + \left(r - \frac{\sigma^2}{2}\right) \frac{\partial u}{\partial x} - r u, & x \in \mathbb{R}, t > 0, \\ u(0, x) = f(x), & x \in \mathbb{R}. \end{cases}\quad (3.5)$$

The following analogue of Theorem 1 now follows.

Theorem 2. *Let $s > 1, \tau \geq 0$. Let σ, r be given positive constants. Let $T = \{T(t) : t \geq 0\}$ be the semigroup governing (3.5) on $X_{s,\tau}$. Then T is a chaotic (C_0) semigroup. The same conclusions hold for the restriction of T to*

$$X_{s,\tau}^{\mathbb{R}} = \{f \in X_{s,\tau} : f \text{ is a real function}\}.$$

The following Example concerning a convection-diffusion type equation can be found in [11, Example 4.12].

Example 1. The operator A defined in the following linear evolution problem

$$\begin{cases} \frac{\partial u}{\partial t} = Au := a \frac{\partial^2 u}{\partial x^2} + b \frac{\partial u}{\partial x} + cu, & x \in [0, \infty), t \geq 0, \\ u(0, t) = 0, & \text{for } t \geq 0, \\ u(x, 0) = f(x) & \text{for } x > 0, \end{cases}$$

in the Hilbert space $L^2(\mathbb{R}_+, \mathbb{C})$ with domain $D(A) = \{f \in W^{2,2}(\mathbb{R}_+) : f(0) = 0\}$ generates an analytic chaotic semigroup, provided that $(a, b, c) \in (0, \infty)^3$ and $0 < c < b^2/2a < 1$.

As a comparison our result treats also the cases $\{b = 0 \text{ and } c < 0\}$ and $\{c < 0 < a \text{ and } a + b + c = 0\}$.

References

- [1] G. Herzog. On the universality of the heat equation. *Math. Nachr.*, 188:169–171, 1997.
- [2] K. G. Grosse-Erdmann. Universal families and hypercyclic operators. *Bull. Amer. Math. Soc.*, 36:345–381, 1999.
- [3] F. Astengo and B. Di Blasio. Dynamics of the heat semigroup in Jacobi analysis. *J. Math. Anal. Appl.*, 391:48–56, 2012.
- [4] H. Emamirad, G. R. Goldstein, and J. A. Goldstein. Chaotic solution for the Black-Scholes equation. *Proc. Amer. Math. Soc.*, 140:2043–2052, 2012.
- [5] H. Emamirad, G. R. Goldstein, and J. A. Goldstein. Corrigendum and improvement for *chaotic solution for the Black-Scholes equation*. *Proc. Amer. Math. Soc.* to appear.
- [6] R. L. Devaney. *An Introduction to Chaotic Dynamical Systems*. Addison-Wesley, New York, 1989.
- [7] N. P. Romanoff. On one parameter operator groups of linear transformation I. *Ann. Math.*, 48:216–233, 1947.
- [8] J. A. Goldstein. *Semigroups of Linear Operators and Applications*. Oxford University Press, New York, 1985.
- [9] J. A. Goldstein, R. M. Mininni, and S. Romanelli. A new explicit formula for the Black-Merton-Scholes equation. In A. Sengupta and P. Sundar, editors, *Infinite Dimensional Stochastic Analysis*, pages 226–235, 2008.
- [10] B. Oksendal. *Stochastic Differential Equation*. Springer-Verlag, Berlin, New York, 1995.
- [11] W. Desch, W. Schappacher, and G. F. Webb. Hypercyclic and chaotic semigroup of linear operator. *Ergodic Theory Dyn. Syst.*, 17:793–819, 1997.