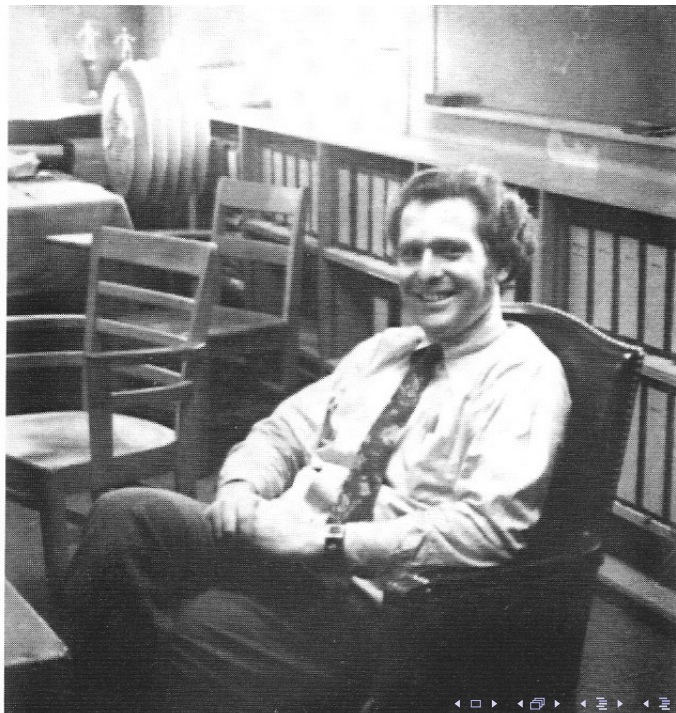


Global solutions
for some reaction-diffusion systems:
 L^∞ , L^p , L^1 and L^2 -strategies

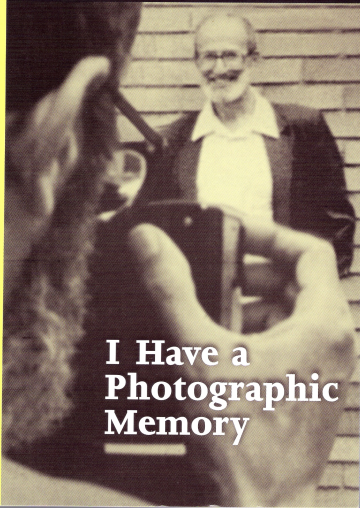
Michel Pierre

Ecole Normale Supérieure de Cachan-Bretagne and Institut de Recherche
Mathématique de Rennes, France

Conference in honor of Jerry Goldstein
Poitiers, June 10th, 2011



PAUL R. HALMOS



**I Have a
Photographic
Memory**

 **AMS**
AMERICAN MATHEMATICAL SOCIETY



An easy O.D.E.

$$\begin{cases} u' = -u v^\beta, \\ v' = u v^\beta \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0, \\ u_0, v_0 \text{ given in } [0, \infty), \end{cases}$$

where $u, v : [0, T) \rightarrow \mathbf{R}$ are the unknown functions. Here $\beta \geq 1$. Local existence of a **nonnegative** unique solution on a maximal interval $[0, T^*)$ is well-known due to the C^1 -property of $(u, v) \rightarrow uv^\beta$. Moreover $u \geq 0, v \geq 0$ and

$$(u + v)'(t) = 0 \Rightarrow (u + v)(t) = u_0 + v_0,$$

so that: $\sup_{t \in [0, T^*)} |u(t)| + |v(t)| < +\infty$,
and therefore

$$T^* = +\infty$$

What happens when diffusion is added?

$$\begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta & \text{in } Q_T = (0, T) \times \Omega \\ \partial_t v - d_2 \Delta v = uv^\beta & \text{in } Q_T = (0, T) \times \Omega \\ \partial_\nu u = \partial_\nu v = 0 & \text{on } \Sigma_T = (0, T) \times \partial\Omega, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

Here $\Omega \subset \mathbf{R}^N$, regular. The total mass is preserved:

$$\int_{\Omega} \partial_t(u + v) - \int_{\Omega} \Delta(d_1 u + d_2 v) = 0.$$

$$\partial_\nu(d_1 u + d_2 v) = 0 \text{ on } \partial\Omega \Rightarrow \int_{\Omega} \Delta(d_1 u + d_2 v) = 0.$$

$$\int_{\Omega} (u + v)(t) = \int_{\Omega} u_0 + v_0$$

Insufficient for global existence!

Local existence for reaction-diffusion systems with L^∞ -data

$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

Theorem: Let $u_0, v_0 \in L^\infty(\Omega)$, $u_0 \geq 0, v_0 \geq 0$. Then, there exist a maximum time $T^* > 0$ and (u, v) unique **classical** nonnegative solution of (S) on $[0, T^*[$. Moreover,

$$\sup_{t \in [0, T^*[} \{ \|u(t)\|_{L^\infty(\Omega)} + \|v(t)\|_{L^\infty(\Omega)} \} < +\infty \Rightarrow [T^* + \infty].$$

Remark: here $\int_\Omega (u + v)(t) = \int_\Omega u_0 + v_0$, that is

$$\sup_{t \in [0, T^*[} \{ \|u(t)\|_{L^1(\Omega)} + \|v(t)\|_{L^1(\Omega)} \} \leq \|u_0\|_{L^1(\Omega)} + \|v_0\|_{L^1(\Omega)}.$$

How does this estimate help for global existence? Very frequent situation in applications !

Same question for the general family of systems:

$$\begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

$d_i > 0$, $f_i : [0, \infty)^m \rightarrow \mathbf{R}$ of class C^1 where

- ▶ **(P)**: Positivity (nonnegativity) is preserved

Same question for the general family of systems:

$$\begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

$d_i > 0$, $f_i : [0, \infty)^m \rightarrow \mathbf{R}$ of class C^1 where

- ▶ **(P)**: Positivity (nonnegativity) is preserved
- ▶ **(M)**: $\sum_{1 \leq i \leq m} f_i \leq 0$ or more generally

Same question for the general family of systems:

$$\begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

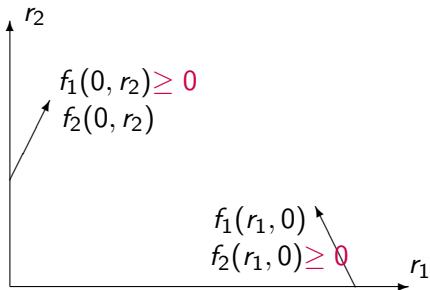
$d_i > 0$, $f_i : [0, \infty)^m \rightarrow \mathbf{R}$ of class C^1 where

- ▶ **(P)**: Positivity (nonnegativity) is preserved
- ▶ **(M)**: $\sum_{1 \leq i \leq m} f_i \leq 0$ or more generally
- ▶ **(M')** $\forall r \in [0, \infty[^m$, $\sum_{1 \leq i \leq m} a_i f_i(r) \leq C[1 + \sum_{1 \leq i \leq m} r_i]$
for some $a_i > 0$

$$(E) \begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

► **(P) Preservation of Positivity:** $\forall i = 1, \dots, m$

$\forall r = (r_1, \dots, r_m) \in [0, \infty[^m, f_i(r_1, \dots, r_{i-1}, 0, r_{i+1}, \dots, r_m) \geq 0.$



$$\begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

- ▶ **(P) Preservation of Positivity:** $\forall i = 1, \dots, m$
 $\forall r \in [0, +\infty[^m, f_i(r_1, \dots, r_{i-1}, 0, r_{i+1}, \dots, r_m) \geq 0.$
- ▶ **(M):** $\sum_{1 \leq i \leq m} f_i(r_1, \dots, r_m) \leq 0 \Rightarrow$ **'Control of the Total Mass':**

$$\forall t \geq 0, \int_{\Omega} \sum_{1 \leq i \leq r} u_i(t, x) dx \leq \int_{\Omega} \sum_{1 \leq i \leq r} u_i^0(x) dx.$$

Add up, integrate on Ω , use $\int_{\Omega} \Delta u_i = \int_{\partial\Omega} \partial_\nu u_i = 0$:

$$\int_{\Omega} \partial_t [\sum u_i(t)] dx = \int_{\Omega} \sum_i f_i(u) dx \leq 0.$$

$$\begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

- ▶ **(P) Preservation of Positivity:** $\forall i = 1, \dots, m$
 $\forall r \in [0, +\infty[^m, f_i(r_1, \dots, r_{i-1}, 0, r_{i+1}, \dots, r_m) \geq 0.$
- ▶ **(M):** $\sum_{1 \leq i \leq m} f_i(r_1, \dots, r_m) \leq 0 \Rightarrow$ **'Control of the Total Mass':**

$$\forall t \geq 0, \int_{\Omega} \sum_{1 \leq i \leq r} u_i(t, x) dx \leq \int_{\Omega} \sum_{1 \leq i \leq r} u_i^0(x) dx.$$

Add up, integrate on Ω , use $\int_{\Omega} \Delta u_i = \int_{\partial\Omega} \partial_\nu u_i = 0$:

$$\int_{\Omega} \partial_t [\sum u_i(t)] dx = \int_{\Omega} \sum_i f_i(u) dx \leq 0.$$

- ▶ $\Rightarrow L^1(\Omega)$ - a priori estimates, uniform in time.

$$\begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

- ▶ **(P) Preservation of Positivity:** $\forall i = 1, \dots, m$
 $\forall r \in [0, +\infty[^m, f_i(r_1, \dots, r_{i-1}, 0, r_{i+1}, \dots, r_m) \geq 0.$
- ▶ **(M):** $\sum_{1 \leq i \leq m} f_i(r_1, \dots, r_m) \leq 0 \Rightarrow$ **'Control of the Total Mass':**

$$\forall t \geq 0, \int_{\Omega} \sum_{1 \leq i \leq r} u_i(t, x) dx \leq \int_{\Omega} \sum_{1 \leq i \leq r} u_i^0(x) dx.$$

Add up, integrate on Ω , use $\int_{\Omega} \Delta u_i = \int_{\partial\Omega} \partial_\nu u_i = 0$:

$$\int_{\Omega} \partial_t [\sum u_i(t)] dx = \int_{\Omega} \sum_i f_i(u) dx \leq 0.$$

- ▶ $\Rightarrow L^1(\Omega)$ - a priori estimates, uniform in time.
- ▶ Remark: same with **(M')**

QUESTION:

What about Global Existence of solutions

under assumption $(\mathbf{P})+(\mathbf{M})??$

or more generally $(\mathbf{P})+(\mathbf{M}') ??$

Explicit examples with property (P)+(M) ou (M')

"Chemical morphogenetic process ("Brusselator", Prigogine)

$$\left\{ \begin{array}{l} \partial_t u - d_1 \Delta u = -uv^2 + b v \\ \partial_t v - d_2 \Delta v = uv^2 - (b + 1) v + a \\ u|_{\partial\Omega} = b/a, \quad v|_{\partial\Omega} = a, \\ a, b, d_1, d_2 > 0. \end{array} \right.$$

Explicit examples with property **(P)**+**(M)**

- ▶ **"Chemical morphogenetic process ("Brusselator", Prigogine)**

$$\begin{cases} \partial_t u - d_1 \Delta u = -uv^2 + bv \\ \partial_t v - d_2 \Delta v = uv^2 - (b+1)v + a \end{cases}$$

See also: Glycolysis model–Gray-Scott models

Explicit examples with property **(P)**+**(M)**

- ▶ "Chemical morphogenetic process ("Brusselator", Prigogine)

$$\begin{cases} \partial_t u - d_1 \Delta u = -uv^2 + bv \\ \partial_t v - d_2 \Delta v = uv^2 - (b+1)v + a \end{cases}$$

See also: Glycolysis model–Gray-Scott models

- ▶ Exothermic combustion in a gas

$$\begin{cases} \partial_t Y - \mu \Delta Y = -H(Y, T) \\ \partial_t T - \lambda \Delta T = q H(Y, T), \end{cases}$$

Y = concentration of a reactant, T = temperature,

Explicit examples with **(P)+(M)**

► Lotka-Volterra Systems (Leung,...)

$$\forall i = 1 \dots m, \quad \partial_t u_i - d_i \Delta u_i = e_i u_i + u_i \sum_{1 \leq j \leq m} p_{ij} u_j,$$

with $e_i, p_{ij} \in \mathbf{R}$ such that for some $a_i > 0$.

$$\forall w \in \mathbf{R}^m, \quad \sum_{i,j=1}^m a_i w_i p_{ij} w_j \leq 0,$$

Explicit examples with **(P)+(M)**

► Lotka-Volterra Systems (Leung,...)

$$\forall i = 1 \dots m, \quad \partial_t u_i - d_i \Delta u_i = e_i u_i + u_i \sum_{1 \leq j \leq m} p_{ij} u_j,$$

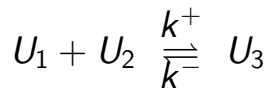
with $e_i, p_{ij} \in \mathbf{R}$ such that for some $a_i > 0$.

$$\forall w \in \mathbf{R}^m, \quad \sum_{i,j=1}^m a_i w_i p_{ij} w_j \leq 0,$$

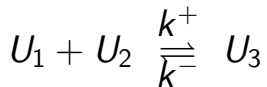
► Model of diffusive calcium dynamics: H.G. Othmer

$$\begin{cases} \partial_t u_1 = d_1 \Delta u_1 + \lambda(\gamma_0 + \gamma_1 u_4)(1 - u_1) - \frac{p_1 u_1^4}{p_2^4 + u_1^4} \\ \partial_t u_2 - d_2 \Delta u_2 = -k_1 u_2 + k'_1 u_3 \\ \partial_t u_3 - d_3 \Delta u_3 = -k'_1 u_3 - k_2 u_1 u_3 + k_1 u_2 + k'_2 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k_2 u_1 u_3 + k'_3 u_5 - k'_2 u_4 - k_3 u_1 u_4 \\ \partial_t u_5 - d_5 \Delta u_5 = k_3 u_1 u_4 - k'_3 u_5. \end{cases}$$

Elementary chemical reactions: a simple example

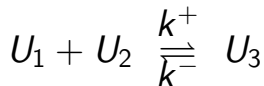


Elementary chemical reactions: a simple example



- ▶ $u_i = u_i(t, x)$ = concentration of U_i

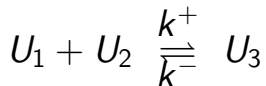
Elementary chemical reactions: a simple example



- ▶ $u_i = u_i(t, x)$ = concentration of U_i
- ▶ *Law of Mass Action*: The instantaneous variation of concentration of each u_i is proportional to the concentration of the reactants:

$$\partial_t u_1 + \nabla \cdot (u_1 \mathbf{V}_1) = k^- u_3 - k^+ u_1 u_2 \quad (\mathbf{V}_1 = \text{velocity})$$

Elementary chemical reactions: a simple example

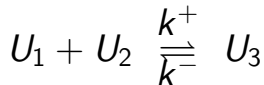


- ▶ $u_i = u_i(t, x)$ = concentration of U_i
- ▶ *Law of Mass Action*: The instantaneous variation of concentration of each u_i is proportional to the concentration of the reactants:

$$\partial_t u_1 + \nabla \cdot (u_1 \mathbf{V}_1) = k^- u_3 - k^+ u_1 u_2 \quad (\mathbf{V}_1 = \text{velocity})$$

- ▶ Fick's diffusion law: $\mathbf{V}_1 = -d_1 \nabla u_1 \Rightarrow \nabla \cdot (u_1 \mathbf{V}_1) = -d_1 \Delta u_1$

Elementary chemical reactions: a simple example



- ▶ $u_i = u_i(t, x)$ = concentration of U_i
- ▶ *Law of Mass Action*: The instantaneous variation of concentration of each u_i is proportional to the concentration of the reactants:

$$\partial_t u_1 + \nabla \cdot (u_1 \mathbf{V}_1) = k^- u_3 - k^+ u_1 u_2 \quad (\mathbf{V}_1 = \text{velocity})$$

- ▶ Fick's diffusion law: $\mathbf{V}_1 = -d_1 \nabla u_1 \Rightarrow \nabla \cdot (u_1 \mathbf{V}_1) = -d_1 \Delta u_1$



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 \end{cases}$$

Note : $f_1 + f_2 + 2f_3 = 0$ and positivity is preserved.

A quadratic model



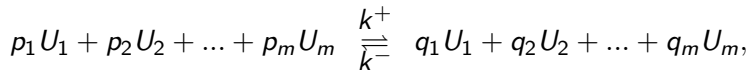
$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

Note: $f_1 + f_2 + f_3 + f_4 = 0$ and positivity is preserved.

L. Desvillettes, K. Fellner, M.P., J. Vovelle, Th. Goudon, A. Vasseur

Superquadratic reaction-diffusion systems.

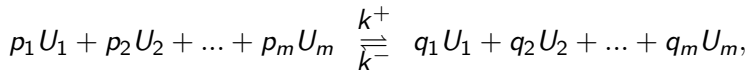
- ▶ A general chemical reaction:



p_i, q_i = nonnegative integers.

Superquadratic reaction-diffusion systems.

- ▶ A general chemical reaction:



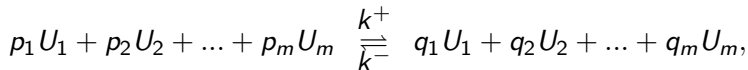
$p_i, q_i =$ nonnegative integers.

- ▶

$$\partial_t u_i - d_i \Delta u_i = (p_i - q_i) \left(k^- \prod_{j=1}^m u_j^{q_j} - k^+ \prod_{j=1}^m u_j^{p_j} \right), \forall i = 1 \dots m.$$

Superquadratic reaction-diffusion systems.

- ▶ A general chemical reaction:



$p_i, q_i = \text{nonnegative integers.}$

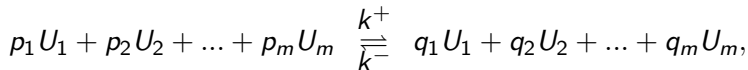


$$\partial_t u_i - d_i \Delta u_i = (p_i - q_i) \left(k^- \prod_{j=1}^m u_j^{q_j} - k^+ \prod_{j=1}^m u_j^{p_j} \right), \forall i = 1 \dots m.$$

- ▶ Here $\sum_i m_i p_i = \sum_i m_i q_i$ for some $m_i \in (0, \infty), i = 1 \dots m$.
This implies (\mathbf{M}') : $\sum_{i=1}^m m_i f_i = 0$.

Superquadratic reaction-diffusion systems.

- ▶ A general chemical reaction:



$p_i, q_i =$ nonnegative integers.



$$\partial_t u_i - d_i \Delta u_i = (p_i - q_i) \left(k^- \prod_{j=1}^m u_j^{q_j} - k^+ \prod_{j=1}^m u_j^{p_j} \right), \forall i = 1 \dots m.$$

- ▶ Here $\sum_i m_i p_i = \sum_i m_i q_i$ for some $m_i \in (0, \infty), i = 1 \dots m$.
This implies (\mathbf{M}') : $\sum_{i=1}^m m_i f_i = 0$.
- ▶ **Global existence completely open in this general situation!**

More examples

► **Diffusive epidemic models: SIR**

S =Susceptibles= can be infected

I =Infectives=infected and transmit disease

R =Removed=immune; $P = S + I + R$

$$\begin{cases} S_t - \nabla \cdot d_1(x)\nabla S = bP - (m + kP)S - g(S, I) \\ I_t - \nabla \cdot d_2(x)\nabla I = -(m + kP)I + g(S, I) - \lambda I \\ R_t - \nabla \cdot d_3(x)\nabla R = -(m + kP)R + \lambda I \end{cases}$$

May be coupled with an extra variable: $S = S(t, x, \text{age})$

More examples

► Diffusive epidemic models: SIR

S =Susceptibles= can be infected

I =Infectives=infected and transmit disease

R =Removed=immune; $P = S + I + R$

$$\begin{cases} S_t - \nabla \cdot d_1(x)\nabla S = bP - (m + kP)S - g(S, I) \\ I_t - \nabla \cdot d_2(x)\nabla I = -(m + kP)I + g(S, I) - \lambda I \\ R_t - \nabla \cdot d_3(x)\nabla R = -(m + kP)R + \lambda I \end{cases}$$

May be coupled with an extra variable: $S = S(t, x, \text{age})\dots$

► Nickel-Iron alloy electrodeposition : N. Alaa, A. Tounsi et al. : $\forall i = 1\dots 5$

$$\begin{cases} \partial_t w_i - d_i(w_i)_{xx} + b(x)(w_i)_x - m_i[w_i\Phi_x]_x = S_i(w) \\ S_1 = S_2 = 0, S_3(w) = S_4(w) = -S_5(w) \\ -[\Phi]_{xx} = \sum_{i=1}^5 z_i w_i, \quad z_i \in R, \quad +\text{bdy cond.} \end{cases}$$

Explicit examples

- ▶ **Modelization of pollutants transfer in atmosphere**
($N = 3$): R. Texier-Picard, MP, (degenerate diffusion):

$$\begin{cases} \partial_t \phi_i = d_i \partial_{zz}^2 \phi_i + \omega \cdot \nabla \phi_i + f_i(\phi) + g_i, \quad \forall i = 1 \dots 20, \\ + \text{Bdy and initial conditions} \end{cases}$$

Explicit examples

- **Modelization of pollutants transfer in atmosphere**
($N = 3$): R. Texier-Picard, MP, (degenerate diffusion):

$$\left\{ \begin{array}{l} \partial_t \phi_i = d_i \partial_{zz}^2 \phi_i + \omega \cdot \nabla \phi_i + f_i(\phi) + g_i, \quad \forall i = 1 \dots 20, \\ + \text{Bdy and initial conditions} \end{array} \right.$$

- **The reaction terms:**

$$\left\{ \begin{array}{l} f_1(\phi) = -k_1 \phi_1 + k_{22} \phi_{19} + k_{25} \phi_{20} + k_{11} \phi_{13} + k_9 \phi_{11} \phi_2 + k_3 \phi_5 \phi_2 \\ \quad + k_2 \phi_2 \phi_4 - k_{23} \phi_1 \phi_4 - k_{14} \phi_1 \phi_6 + k_{12} \phi_{10} \phi_2 - k_{10} \phi_{11} \phi_1 - k_{24} \phi_{19} \phi_1, \\ f_2(\phi) = k_1 \phi_1 + k_{21} \phi_{19} - k_9 \phi_{11} \phi_2 - k_3 \phi_5 \phi_2 - k_2 \phi_2 \phi_4 - k_{12} \phi_{10} \phi_2 \\ f_3(\phi) = k_1 \phi_1 + k_{17} \phi_4 + k_{19} \phi_{16} + k_{22} \phi_{19} - k_{15} \phi_3 \\ f_4(\phi) = -k_{17} \phi_4 + k_{15} \phi_3 - k_{16} \phi_4 - k_2 \phi_2 \phi_4 - k_{23} \phi_1 \phi_4 \\ f_5(\phi) = 2k_4 \phi_7 + k_7 \phi_9 + k_{13} \phi_{14} + k_6 \phi_7 \phi_6 - k_3 \phi_5 \phi_2 + k_{20} \phi_{17} \phi_6 \\ f_6(\phi) = 2k_{18} \phi_{16} - k_8 \phi_9 \phi_6 - k_6 \phi_7 \phi_6 + k_3 \phi_5 \phi_2 - k_{20} \phi_{17} \phi_6 - k_{14} \phi_1 \phi_6 \\ f_7(\phi) = -k_4 \phi_7 - k_5 \phi_7 + k_{13} \phi_{14} - k_6 \phi_7 \phi_6 \\ f_8(\phi) = k_4 \phi_7 + k_5 \phi_7 + k_7 \phi_9 + k_6 \phi_7 \phi_6 \\ f_9(\phi) = -k_7 \phi_9 - k_8 \phi_9 \phi_6 \\ f_{10}(\phi) = k_7 \phi_9 + k_9 \phi_{11} \phi_2 - k_{12} \phi_{10} \phi_2 \\ f_{11}(\phi) = k_{11} \phi_{13} - k_9 \phi_{11} \phi_2 + k_8 \phi_9 \phi_6 - k_{10} \phi_{11} \phi_1 \\ f_{12}(\phi) = k_9 \phi_{11} \phi_2 \\ f_{13}(\phi) = -k_{11} \phi_{13} + k_{10} \phi_{11} \phi_1 \\ f_{14}(\phi) = -k_{13} \phi_{14} + k_{12} \phi_{10} \phi_2 \\ f_{15}(\phi) = k_{14} \phi_1 \phi_6 \\ f_{16}(\phi) = -k_{19} \phi_{16} - k_{18} \phi_{16} + k_{16} \phi_4 \\ f_{17}(\phi) = -k_{20} \phi_{17} \phi_6 \\ f_{18}(\phi) = k_{20} \phi_{17} \phi_6 \\ f_{19}(\phi) = -k_{21} \phi_{19} - k_{22} \phi_{19} + k_{25} \phi_{20} + k_{23} \phi_1 \phi_4 - k_{24} \phi_{19} \phi_1 \\ f_{20}(\phi) = -k_{25} \phi_{20} + k_{24} \phi_{19} \phi_1. \end{array} \right.$$

Back to the model example: what about L^∞ -estimates?



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

By maximum principle

$$\partial_t u - d_1 \Delta u \leq 0 \Rightarrow \|u(t)\|_{L^\infty(\Omega)} \leq \|u_0\|_{L^\infty(\Omega)}$$

Back to the model example: what about L^∞ -estimates?



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

By maximum principle

$$\partial_t u - d_1 \Delta u \leq 0 \Rightarrow \|u(t)\|_{L^\infty(\Omega)} \leq \|u_0\|_{L^\infty(\Omega)}$$

► If $d_1 = d_2 = d$: $\partial_t(u + v) - d\Delta(u + v) = 0$

$$\Rightarrow \|u(t) + v(t)\|_{L^\infty(\Omega)} \leq \|u_0 + v_0\|_{L^\infty(\Omega)}$$

By Positivity, $\forall t \in [0, T^*)$,

$$\Rightarrow \|u(t)\|_{L^\infty(\Omega)}, \|v(t)\|_{L^\infty(\Omega)} \leq \|u_0 + v_0\|_{L^\infty(\Omega)}$$

$$\Rightarrow T^* = +\infty.$$

Back to the model example: what about L^∞ -estimates?



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

By maximum principle

$$\partial_t u - d_1 \Delta u \leq 0 \Rightarrow \|u(t)\|_{L^\infty(\Omega)} \leq \|u_0\|_{L^\infty(\Omega)}$$

▶ If $d_1 = d_2 = d$: $\partial_t(u + v) - d\Delta(u + v) = 0$

$$\Rightarrow \|u(t) + v(t)\|_{L^\infty(\Omega)} \leq \|u_0 + v_0\|_{L^\infty(\Omega)}$$

By Positivity, $\forall t \in [0, T^*)$,

$$\Rightarrow \|u(t)\|_{L^\infty(\Omega)}, \|v(t)\|_{L^\infty(\Omega)} \leq \|u_0 + v_0\|_{L^\infty(\Omega)}$$

$$\Rightarrow T^* = +\infty.$$

▶ What happens when $d_1 \neq d_2$?

A general L^p -approach



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

$$\partial_t v - d_2 \Delta v = -[\partial_t u - d_1 \Delta u], \quad u \in L^\infty(Q_{T^*}).$$

$$v = -[\partial_t - d_2 \Delta]^{-1} (\partial_t - d_1 \Delta) u (= \mathcal{A}u).$$

A general L^p -approach



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

$$\partial_t v - d_2 \Delta v = -[\partial_t u - d_1 \Delta u], \quad u \in L^\infty(Q_{T^*}).$$

$$v = -[\partial_t - d_2 \Delta]^{-1} (\partial_t - d_1 \Delta) u (= \mathcal{A}u).$$

- **Lemma:** the operator \mathcal{A} is continuous from $L^p(Q_T)$ into $L^p(Q_T)$ for all $p \in]1, \infty[$ and all $T > 0$.
[dual statement of the L^p -regularity theory for the heat operator]

$$\Rightarrow \forall p < +\infty, \quad \|v\|_{L^p(Q_{T^*})} < +\infty$$

A general L^p -approach



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

$$\partial_t v - d_2 \Delta v = -[\partial_t u - d_1 \Delta u], \quad u \in L^\infty(Q_{T^*}).$$

$$v = -[\partial_t - d_2 \Delta]^{-1} (\partial_t - d_1 \Delta) u (= \mathcal{A}u).$$

- ▶ **Lemma:** the operator \mathcal{A} is continuous from $L^p(Q_T)$ into $L^p(Q_T)$ for all $p \in]1, \infty[$ and all $T > 0$.
[dual statement of the L^p -regularity theory for the heat operator]

$$\Rightarrow \forall p < +\infty, \quad \|v\|_{L^p(Q_{T^*})} < +\infty$$

- ▶ Next: $\|v\|_{L^\infty(Q_{T^*})} \leq C \|uv^\beta\|_{L^q(Q_{T^*})}$ si $q > (d+1)/2$

A general L^p -approach



$$(S) \begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ on } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ on } Q_T \\ \partial_\nu u = \partial_\nu v = 0 \text{ on } \Sigma_T, \\ u(0) = u_0 \geq 0, \quad v(0) = v_0 \geq 0. \end{cases}$$

$$\partial_t v - d_2 \Delta v = -[\partial_t u - d_1 \Delta u], \quad u \in L^\infty(Q_{T^*}).$$

$$v = -[\partial_t - d_2 \Delta]^{-1} (\partial_t - d_1 \Delta) u (= \mathcal{A}u).$$

- ▶ **Lemma:** the operator \mathcal{A} is continuous from $L^p(Q_T)$ into $L^p(Q_T)$ for all $p \in]1, \infty[$ and all $T > 0$.
[dual statement of the L^p -regularity theory for the heat operator]

$$\Rightarrow \forall p < +\infty, \quad \|v\|_{L^p(Q_{T^*})} < +\infty$$

- ▶ Next: $\|v\|_{L^\infty(Q_{T^*})} \leq C \|uv^\beta\|_{L^q(Q_{T^*})}$ si $q > (d+1)/2$
- ▶ Therefore $\|v\|_{L^\infty(Q_{T^*})} < +\infty$ et $T^* = +\infty$.

Extensions and limits of the L^p -approach

- ▶ The same approach provides global existence
 - for the "Brusselator", for the epidemic models SIR
 - for the 3×3 system

$$U_1 + U_2 \xrightleftharpoons[k^-]{k^+} U_3 : \begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 \end{cases}$$

Extensions and limits of the L^p -approach

- ▶ The same approach provides global existence
 - for the "Brusselator", for the epidemic models SIR
 - for the 3×3 system

$$U_1 + U_2 \stackrel{k^+}{\rightleftharpoons} U_3 : \begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 \end{cases}$$

- ▶ More generally it applies to $m \times m$ systems if there exists a **triangular** invertible matrix Q with nonnegative entries such that

$$\forall r \in [0, \infty)^m, Q f(r) \leq [1 + \sum_{1 \leq i \leq m} r_i] \mathbf{b},$$

for some $\mathbf{b} \in \mathbf{R}^m, f = (f_1, \dots, f_m)^t$.

Extensions and limits of the L^p -approach

- ▶ L^p -approach does not apply to

$$\begin{cases} \partial_t u - d_1 \Delta u = -ue^{v^2} & \text{in } Q_T \\ \partial_t v - d_2 \Delta v = ue^{v^2} & \text{in } Q_T \end{cases}$$

Extensions and limits of the L^p -approach

- ▶ L^p -approach does not apply to

$$\begin{cases} \partial_t u - d_1 \Delta u = -ue^{v^2} & \text{in } Q_T \\ \partial_t v - d_2 \Delta v = ue^{v^2} & \text{in } Q_T \end{cases}$$

- ▶ neither to the system

$$\begin{cases} \partial_t u - d_1 \Delta u = u^3 v^2 - u^2 v^3 & \text{in } Q_T \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 & \text{in } Q_T \end{cases}$$

Extensions and limits of the L^p -approach

- ▶ L^p -approach does not apply to

$$\begin{cases} \partial_t u - d_1 \Delta u = -ue^{v^2} & \text{in } Q_T \\ \partial_t v - d_2 \Delta v = ue^{v^2} & \text{in } Q_T \end{cases}$$

- ▶ neither to the system

$$\begin{cases} \partial_t u - d_1 \Delta u = u^3 v^2 - u^2 v^3 & \text{in } Q_T \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 & \text{in } Q_T \end{cases}$$

- ▶ and even not to the "better" system with $\lambda \in [0, 1[$

$$\begin{cases} \partial_t u - d_1 \Delta u = \lambda u^3 v^2 - u^2 v^3 & \text{in } Q_T, \\ \partial_t v - d_2 \Delta v = u^2 v^3 - \lambda u^3 v^2 & \text{in } Q_T \end{cases}$$

where $f(u, v) + g(u, v) \leq 0$ et also $f(u, v) + \lambda g(u, v) \leq 0$

Finite time L^∞ -blow up may appear!



$$\begin{cases} \partial_t u - d_1 \Delta u = f(u, v) \text{ in } Q_T, \\ \partial_t v - d_2 \Delta v = g(u, v) \text{ in } Q_T \end{cases}$$

Theorem: (D. Schmitt, MP) One can find polynomial nonlinearities f, g satisfying **(P)** and

$$\mathbf{(M)} \quad f + g \leq 0, \text{ and also } : \exists \lambda \in [0, 1[, f + \lambda g \leq 0,$$

and for which there exists $T^* < +\infty$ with

$$\lim_{t \rightarrow T^*} \|u(t)\|_{L^\infty(\Omega)} = \lim_{t \rightarrow T^*} \|v(t)\|_{L^\infty(\Omega)} = +\infty.$$

Finite time L^∞ -blow up may appear!



$$\begin{cases} \partial_t u - d_1 \Delta u = f(u, v) \text{ in } Q_T, \\ \partial_t v - d_2 \Delta v = g(u, v) \text{ in } Q_T \end{cases}$$

Theorem: (D. Schmitt, MP) One can find polynomial nonlinearities f, g satisfying **(P)** and

$$\textbf{(M)} \quad f + g \leq 0, \text{ and also : } \exists \lambda \in [0, 1[, f + \lambda g \leq 0,$$

and for which there exists $T^* < +\infty$ with

$$\lim_{t \rightarrow T^*} \|u(t)\|_{L^\infty(\Omega)} = \lim_{t \rightarrow T^*} \|v(t)\|_{L^\infty(\Omega)} = +\infty.$$

- ▶ The blow up is similar to $u(t, x) = \frac{1}{(T^* - t)^2 + |x|^2}$ The solution goes out of $L^\infty(\Omega)$ at $t = T^*$, but still exists for $t > T^*$. — — — > Incomplete blow up !

CONCLUSION at this stage:

Look rather for *weak solutions* which are allowed to go out of $L^\infty(\Omega)$ from time to time or even often.

We ask the nonlinearities to be at least in $L^1(Q_T)$.

$$f_i(u) \in L^1(Q_T) ?$$

An L^1 -approach

$$(S) \begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

L^1 -Theorem. Assume the two conditions **(P)**+ **(M')** hold. Assume moreover that the following a priori estimate holds:

$$\forall i = 1, \dots, m, \quad \int_{Q_T} |f_i(u)| \leq C.$$

Then, there exists a **global weak solution** for System (S).

L^1 -Theorem applies to many situations

$$\begin{cases} \partial_t u - d_1 \Delta u = -ue^{v^2} & \text{in } Q_T \\ \partial_t v - d_2 \Delta v = ue^{v^2} & \text{in } Q_T \end{cases}$$

$$\int_{\Omega} u(T) + \int_{Q_T} ue^{v^2} = \int_{\Omega} u_0,$$

whence the $L^1(Q_T)$ -estimate of the nonlinearity.

L^1 -Theorem applies to many situations



$$\begin{cases} \partial_t u - d_1 \Delta u = \lambda u^3 v^2 - u^2 v^3 \text{ in } Q_T, \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 \text{ in } Q_T \end{cases}$$

L^1 -Theorem applies to many situations



$$\begin{cases} \partial_t u - d_1 \Delta u = \lambda u^3 v^2 - u^2 v^3 & \text{in } Q_T, \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 & \text{in } Q_T \end{cases}$$



$$\int_{\Omega} u(T) + \int_{Q_T} u^2 v^3 = \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0.$$

$$\int_{\Omega} v(T) + \int_{Q_T} u^3 v^2 = \int_{Q_T} u^2 v^3 + \int_{\Omega} v_0$$

L^1 -Theorem applies to many situations



$$\begin{cases} \partial_t u - d_1 \Delta u = \lambda u^3 v^2 - u^2 v^3 & \text{in } Q_T, \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 & \text{in } Q_T \end{cases}$$



$$\int_{\Omega} u(T) + \int_{Q_T} u^2 v^3 = \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0.$$

$$\int_{\Omega} v(T) + \int_{Q_T} u^3 v^2 = \int_{Q_T} u^2 v^3 + \int_{\Omega} v_0$$



$$\Rightarrow \int_{Q_T} u^3 v^2 \leq \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0 + v_0$$

L^1 -Theorem applies to many situations



$$\begin{cases} \partial_t u - d_1 \Delta u = \lambda u^3 v^2 - u^2 v^3 & \text{in } Q_T, \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 & \text{in } Q_T \end{cases}$$



$$\int_{\Omega} u(T) + \int_{Q_T} u^2 v^3 = \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0.$$

$$\int_{\Omega} v(T) + \int_{Q_T} u^3 v^2 = \int_{Q_T} u^2 v^3 + \int_{\Omega} v_0$$



$$\Rightarrow \int_{Q_T} u^3 v^2 \leq \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0 + v_0$$



$$\text{For } \lambda < 1 : \Rightarrow \int_{Q_T} u^3 v^2 < +\infty, \int_{Q_T} u^2 v^3 < +\infty$$

L^1 -Theorem applies to many situations



$$\begin{cases} \partial_t u - d_1 \Delta u = \lambda u^3 v^2 - u^2 v^3 & \text{in } Q_T, \\ \partial_t v - d_2 \Delta v = u^2 v^3 - u^3 v^2 & \text{in } Q_T \end{cases}$$



$$\int_{\Omega} u(T) + \int_{Q_T} u^2 v^3 = \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0.$$

$$\int_{\Omega} v(T) + \int_{Q_T} u^3 v^2 = \int_{Q_T} u^2 v^3 + \int_{\Omega} v_0$$



$$\Rightarrow \int_{Q_T} u^3 v^2 \leq \lambda \int_{Q_T} u^3 v^2 + \int_{\Omega} u_0 + v_0$$



$$\text{For } \lambda < 1 : \Rightarrow \int_{Q_T} u^3 v^2 < +\infty, \int_{Q_T} u^2 v^3 < +\infty$$

► Open problem if $\lambda = 1$: L^1 -estimate of the nonlinearity??

L^1 -Theorem applies to many situations

More generally it applies if there exists an invertible matrix Q with nonnegative entries such that

$$\forall r \in [0, \infty)^m, Qf(r) \leq [1 + \sum_{1 \leq i \leq m} r_i] \mathbf{b},$$

for some $\mathbf{b} \in \mathbf{R}^m, f = (f_1, \dots, f_m)^t$.

A surprising a priori L^2 -estimate for these systems



$$(S) \begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

L^2 -Theorem. Assume **(P)**+**(M')**. Then, the following a priori estimate holds for the solutions of (S):

$$\forall i = 1, \dots, m, \forall T > 0, \int_{Q_T} u_i^2 \leq C.$$

A surprising a priori L^2 -estimate for these systems



$$(S) \begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

L^2 -Theorem. Assume **(P)**+**(M')**. Then, the following a priori estimate holds for the solutions of (S):

$$\forall i = 1, \dots, m, \forall T > 0, \int_{Q_T} u_i^2 \leq C.$$

- ▶ **Corollary of the L^1 - and L^2 -Theorems:** Assume **(P)**,**(M')** and f_i is at most quadratic. Then, System (S) has a global weak solution.

A surprising a priori L^2 -estimate for these systems



$$(S) \begin{cases} \forall i = 1, \dots, m \\ \partial_t u_i - d_i \Delta u_i = f_i(u_1, u_2, \dots, u_m) & \text{in } Q_T \\ \partial_\nu u_i = 0 & \text{on } \Sigma_T \\ u_i(0, \cdot) = u_i^0(\cdot) \geq 0. \end{cases}$$

L^2 -Theorem. Assume **(P)**+**(M')**. Then, the following a priori estimate holds for the solutions of (S):

$$\forall i = 1, \dots, m, \forall T > 0, \int_{Q_T} u_i^2 \leq C.$$

- ▶ **Corollary of the L^1 - and L^2 -Theorems:** Assume **(P)**,**(M')** and f_i is at most quadratic. Then, System (S) has a global weak solution.
- ▶ Recall that nonlinearities are quadratic in many examples.

Application to the quadratic chemical reaction:



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

Application to the quadratic chemical reaction:



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

- ▶ Global existence of a **weak solution**

Application to the quadratic chemical reaction:



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

- ▶ Global existence of a **weak solution**
- ▶ The L^p -approach does not work

Application to the quadratic chemical reaction:



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

- ▶ Global existence of a **weak solution**
- ▶ The L^p -approach does not work
- ▶ This solution is regular (=classical) in dimension $N = 1, 2$

Application to the quadratic chemical reaction:



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

- ▶ Global existence of a **weak solution**
- ▶ The L^p -approach does not work
- ▶ This solution is regular (=classical) in dimension $N = 1, 2$
- ▶ For $N \geq 3$, the set of points around which the solution is unbounded is "small" in the sense that its Hausdorff dimension is at most $(N^2 - 4)/N$ (Th. Goudon, A. Vasseur)

Application to the quadratic chemical reaction:



$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -k^+ u_1 u_2 + k^- u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = k^+ u_1 u_2 - k^- u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = k^+ u_1 u_2 - k^- u_3 u_4 \end{cases}$$

- ▶ Global existence of a **weak solution**
- ▶ The L^p -approach does not work
- ▶ This solution is regular (=classical) in dimension $N = 1, 2$
- ▶ For $N \geq 3$, the set of points around which the solution is unbounded is "small" in the sense that its Hausdorff dimension is at most $(N^2 - 4)/N$ (Th. Goudon, A. Vasseur)
- ▶ **Open problem: does the solution blow up in $L^\infty(\Omega)$ in finite time or not??**

Idea of the proof of the L^2 -estimate



$$\partial_t \left(\sum_i u_i \right) - \Delta \left(\sum_i d_i u_i \right) \leq 0.$$

Idea of the proof of the L^2 -estimate



$$\partial_t \left(\sum_i u_i \right) - \Delta \left(\sum_i d_i u_i \right) \leq 0.$$



$$\partial_t W - \Delta(aW) \leq 0, \quad W = \sum_i u_i \quad a = \frac{\sum_i d_i u_i}{\sum_i u_i}$$

Idea of the proof of the L^2 -estimate



$$\partial_t \left(\sum_i u_i \right) - \Delta \left(\sum_i d_i u_i \right) \leq 0.$$



$$\partial_t W - \Delta (a W) \leq 0, \quad W = \sum_i u_i \quad a = \frac{\sum_i d_i u_i}{\sum_i u_i}$$



$$0 \leq \min_i d_i \leq a = \frac{\sum_i d_i u_i}{\sum_i u_i} \leq \max_i d_i < +\infty$$

Idea of the proof of the L^2 -estimate



$$\partial_t \left(\sum_i u_i \right) - \Delta \left(\sum_i d_i u_i \right) \leq 0.$$



$$\partial_t W - \Delta(aW) \leq 0, \quad W = \sum_i u_i \quad a = \frac{\sum_i d_i u_i}{\sum_i u_i}$$



$$0 \leq \min_i d_i \leq a = \frac{\sum_i d_i u_i}{\sum_i u_i} \leq \max_i d_i < +\infty$$

- ▶ The operator $W \rightarrow \partial_t W - \Delta(aW)$ is not of divergence form and a is not continuous, but **bounded from above and from below** so that the operator is parabolic and, at least:

$$\|W\|_{L^2(Q_T)} \leq C \|W_0\|_{L^2(\Omega)}.$$

Idea of the proof of the L^2 -estimate



$$\partial_t \left(\sum_i u_i \right) - \Delta \left(\sum_i d_i u_i \right) \leq 0.$$



$$\partial_t W - \Delta(aW) \leq 0, \quad W = \sum_i u_i \quad a = \frac{\sum_i d_i u_i}{\sum_i u_i}$$



$$0 \leq \min_i d_i \leq a = \frac{\sum_i d_i u_i}{\sum_i u_i} \leq \max_i d_i < +\infty$$

- ▶ The operator $W \rightarrow \partial_t W - \Delta(aW)$ is not of divergence form and a is not continuous, but **bounded from above and from below** so that the operator is parabolic and, at least:

$$\|W\|_{L^2(Q_T)} \leq C \|W_0\|_{L^2(\Omega)}.$$

- ▶ We may even show that the mapping $W_0 \in L^2(\Omega) \rightarrow W \in L^2(Q_T)$ is **compact** where $\partial_t W - \Delta(aW) = 0$, $W(0) = W_0$.

An application of the L^2 -compactness



$$U_1 + U_2 \stackrel{\frac{1}{k_1}}{\rightrightarrows} C \stackrel{\frac{k_2}{1}}{\rightrightarrows} U_3 + U_4$$

An application of the L^2 -compactness



- ▶ The intermediate **C** is highly reactive, so that we may assume that $k_1, k_2 \rightarrow +\infty$.

What is the limit kinetics when space diffusion occurs?

An application of the L^2 -compactness



$$U_1 + U_2 \xrightleftharpoons[k_1]{1} C \xrightleftharpoons[1]{k_2} U_3 + U_4$$

- ▶ The intermediate **C** is highly reactive, so that we may assume that $k_1, k_2 \rightarrow +\infty$.

What is the limit kinetics when space diffusion occurs?

- ▶ Mass Action law + Fick's diffusion law lead to the system

$$\left\{ \begin{array}{l} \partial_t u_1 - d_1 \Delta u_1 = -u_1 u_2 + k_1 c \\ \partial_t u_2 - d_2 \Delta u_2 = -u_1 u_2 + k_1 c \\ \partial_t c - d_c \Delta c = u_1 u_2 - (k_1 + k_2) c + u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = -u_3 u_4 + k_2 c \\ \partial_t u_4 - d_4 \Delta u_4 = -u_3 u_4 + k_2 c, \end{array} \right\} \text{ on } Q_T$$

An application of the L^2 -compactness



$$U_1 + U_2 \xrightleftharpoons[k_1]{1} C \xrightleftharpoons[1]{k_2} U_3 + U_4$$

- ▶ The intermediate C is highly reactive, so that we may assume that $k_1, k_2 \rightarrow +\infty$.

What is the limit kinetics when space diffusion occurs?

- ▶ Mass Action law + Fick's diffusion law lead to the system

$$\left\{ \begin{array}{l} \partial_t u_1 - d_1 \Delta u_1 = -u_1 u_2 + k_1 c \\ \partial_t u_2 - d_2 \Delta u_2 = -u_1 u_2 + k_1 c \\ \partial_t c - d_c \Delta c = u_1 u_2 - (k_1 + k_2) c + u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = -u_3 u_4 + k_2 c \\ \partial_t u_4 - d_4 \Delta u_4 = -u_3 u_4 + k_2 c, \end{array} \right\} \text{ on } Q_T$$

- ▶ The L^p -approach applies to this system so that global existence of classical solutions holds!

The limit system

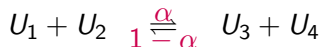
Theorem. The solution $(u_1^k, u_2^k, c^k, u_3^k, u_4^k)$, $k = (k_1, k_2)$ of the previous system converges as $k_1 + k_2 \rightarrow +\infty$ in $L^2(Q_T)^5$ for all $T > 0$ to $(u_1, u_2, 0, u_3, u_4)$ solution of

$$\begin{cases} \partial_t u_1 - d_1 \Delta u_1 = -\alpha u_1 u_2 + \beta u_3 u_4 \\ \partial_t u_2 - d_2 \Delta u_2 = -\alpha u_1 u_2 + \beta u_3 u_4 \\ \partial_t u_3 - d_3 \Delta u_3 = \alpha u_1 u_2 - \beta u_3 u_4 \\ \partial_t u_4 - d_4 \Delta u_4 = \alpha u_1 u_2 - \beta u_3 u_4, \end{cases}$$

where $\alpha = \lim_{k_1+k_2 \rightarrow \infty} \frac{k_2}{k_1+k_2}$, $\beta = 1 - \alpha$. The chemical reaction



"tends" to the limit chemical reaction:



Other boundary conditions

- ▶ All the previous results extends to Dirichlet or Robin type boundary conditions, **assuming they are all of the same type in all equations**

Other boundary conditions

- ▶ All the previous results extends to Dirichlet or Robin type boundary conditions, **assuming they are all of the same type in all equations**
- ▶ Blow up in finite time may occur in the system

$$\begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ in } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ in } Q_T \\ u = 1, \partial_\nu v = 0 \text{ on } \Sigma_T. \end{cases}$$

Other boundary conditions

- ▶ All the previous results extends to Dirichlet or Robin type boundary conditions, **assuming they are all of the same type in all equations**
- ▶ Blow up in finite time may occur in the system

$$\begin{cases} \partial_t u - d_1 \Delta u = -uv^\beta \text{ in } Q_T \\ \partial_t v - d_2 \Delta v = uv^\beta \text{ in } Q_T \\ u = 1, \partial_\nu v = 0 \text{ on } \Sigma_T. \end{cases}$$

- ▶ Work in progress with Jerry, Gisele + M. Meyries (Karlsruhe) on **nonlinear boundary conditions**, including **Wentzell type**, like

$$\begin{cases} \partial_t u_i - d_i \Delta u_i = f_i(u) \text{ in } Q_T \\ \sigma \partial_t u_i + d_i \partial_\nu u_i - \delta_i \Delta_{\partial\Omega} u_i = g_i(u) \text{ on } \Sigma_T \end{cases}$$

with $\sigma, \delta_i \geq 0$.

More open problems

- atmospheric vertical diffusion/transport of pollutants (**partial results with assumptions on the transport**).
- Uniqueness of weak solutions...or more precisely, what is the way to select the "best" solution, since for instance there is even not uniqueness for $\partial_t u - \Delta u = u^3$.
- how to define renormalized solutions for systems???
- what about initial data in $L^1(\Omega)$?
- what about **nonlinear diffusions??**

$$(E) \begin{cases} \partial_t u - \Delta u^m = f(u, v) & \text{on } Q_T \\ \partial_t v - \Delta v^p = g(u, v) & \text{on } Q_T, \end{cases}$$

- Same questions with general Lyapounov type of structure

$$h'_1(u)f(u, v) + h'_2(v)g(u, v) \leq 0, \quad h_1, h_2 \text{ convex.}$$

- same questions for elliptic systems.

